

# Oxygen-vacancy defects in electron-irradiated Si: the role of carbon in their behavior

Londos, C.A. , Sgourou, E.N. and Chroneos, A.

Postprint deposited in [Curve](#) January 2016

**Original citation:**

Londos, C.A. , Sgourou, E.N. and Chroneos, A. (2013) Oxygen-vacancy defects in electron-irradiated Si: the role of carbon in their behavior. Journal of Materials Science: Materials in Electronics, volume 25 (2): 914-921. DOI: 10.1007/s10854-013-1664-6

<http://dx.doi.org/10.1007/s10854-013-1664-6>

Springer US

The final publication is available at Springer via <http://dx.doi.org/10.1007/s10854-013-1664-6>

**Copyright © and Moral Rights are retained by the author(s) and/ or other copyright owners. A copy can be downloaded for personal non-commercial research or study, without prior permission or charge. This item cannot be reproduced or quoted extensively from without first obtaining permission in writing from the copyright holder(s). The content must not be changed in any way or sold commercially in any format or medium without the formal permission of the copyright holders.**

**CURVE is the Institutional Repository for Coventry University**

<http://curve.coventry.ac.uk/open>

# Oxygen-vacancy defects in electron-irradiated Si: the role of carbon in their behavior

C. A. Londos<sup>1</sup>, E. N. Sgourou<sup>1</sup>, A. Chroneos<sup>2,3</sup>

<sup>1</sup>*University of Athens, Solid State Physics Section, Panepistimiopolis Zografos, Athens 157 84, Greece*

<sup>2</sup>*Engineering and Innovation, The Open University, Milton Keynes MK7 6AA, United Kingdom*

<sup>3</sup>*Department of Materials, Imperial College London, London SW7 2AZ, UK*

## Abstract

This is a detailed study of the effect of carbon impurity on the production of the oxygen – vacancy (VO) pair and its conversion to the VO<sub>2</sub> defect, in electron-irradiated Czochralski silicon (Cz-Si) material, by means of infrared spectroscopy. Upon irradiation vacancies are trapped by oxygen atoms to form VO pairs and it was determined that the presence of carbon enhances the production of this pair. This is attributed to the tendency of carbon to capture self-interstitials, thus decreasing the annihilation rate between vacancies and self-interstitials produced during the irradiation and therefore increasing the availability of vacancies to pair with oxygen atoms in the course of irradiation. Upon annealing a number of VO pairs are captured by oxygen atoms to form VO<sub>2</sub> defects. It was determined that the percentage of the VO converted to the VO<sub>2</sub> defects decreases as the concentration of carbon increases. The phenomenon is discussed in terms of the various reaction channels that VO pair participates upon annealing and how the presence of carbon impacts the balance between these reactions, affecting the final products of the involved processes. Finally, an opposite trend for the conversion of the VO<sub>2</sub> to the VO<sub>3</sub> defect was observed. The percentage of VO<sub>2</sub> converted to VO<sub>3</sub> is enhanced with the increase of carbon content. This finding is discussed in terms of the effect of carbon in the oxygen diffusivity and the final impact on the reaction  $VO_2 + O_i \rightarrow VO_3$ .

## 1 Introduction

Oxygen is the dominant impurity incorporated during processing in Czochralski grown Si and the oxygen-related defects have been investigated for decades [1-3]. Oxygen atoms, located at interstitial sites in the Si lattice ( $O_i$ ), are the most effective sinks for the vacancies created in the course of irradiation leading to the formation of the VO pairs ( $V + O_i \rightarrow VO$ ) the well-known A-centers. There are many technological and fundamental research reasons that demand a detailed understanding of the VO defect properties and behavior: Firstly, the defect introduces an acceptor level [4-6] in the forbidden energy gap at around  $E_c - 0.17$  eV. Due to its electrical activity the A-center acts as a recombination center [7,8], resulting in the deterioration of devices especially those operating in radiation environment, for instance Si detectors. Secondly, upon annealing VO interacts with oxygen impurities and vacancies leading to the formation [9] of various  $V_nO_m$  defects. Members of these families of defects cause [10] leakage currents in p-n junctions thus affecting negatively their properties. Thirdly, VO pair has been proposed [11,12] to act as a vehicle for oxygen aggregation processes in Si. Indeed, although the activation [13] energy for oxygen diffusion in Si is 2.53 eV, oxygen aggregation occurs with much lower activation energy [14] of 1.8 eV. To explain this, fast diffusing oxygen containing species are needed and the oxygen-vacancy pair was considered [12] as a potential candidate. Fourthly, on thermal processing of the Si material a number of VO pairs are captured by oxygen atoms to form  $VO_2$  defects. The latter defect exhibits important properties, for instance the center appears to be bistable [15] and most importantly although  $VO_2$  defect is electrically neutral, in its metastable configuration it manifests electrical activity. Precursor structures in the course of its final formation during the VO to the  $VO_2$  conversion have also been reported [16]. More over it has been suggested that the  $VO_2$  defect acts as nuclei for oxygen precipitation [17,18]. Aggregation and precipitation of oxygen are very important processes affecting the quality of Si-based

devices. Thus, any information on the formation of the VO defects and the mechanisms that govern its conversion to the VO<sub>2</sub> defect is considered very useful from the technological point of view, mainly for developing a fabrication process of Si-based devices with enhanced functionality. There is a fifth important reason: the formation and anneal of VO pair involves reactions with intrinsic defects (vacancies and self-interstitials) and **this information can be used** to control defect processes. To this end, the effect of other impurities, such as carbon present in the lattice on the behavior and properties of the VO and the VO<sub>2</sub> defects, is valuable.

Carbon is the second important impurity present in the lattice of as-grown Si [19-23]. It is an isovalent impurity in Si incorporated at substitutional sites in the lattice (C<sub>s</sub>) and it readily traps self-interstitials produced in the course of irradiation. As a result, carbon atoms are injected at interstitial sites (C<sub>i</sub>). Carbon interstitials exhibit in general a very interesting behavior [19,24-26], and being very mobile at room temperature they interact with primary defects and mainly with other C<sub>s</sub> and O<sub>i</sub> atoms to form the C<sub>s</sub>C<sub>i</sub> and C<sub>i</sub>O<sub>i</sub> defects, respectively. The latter defect is the most significant carbon-related defect and it is expected to impact the behavior of VO defect in carbon containing Cz-Si. It is therefore important to understand how carbon affects the mechanisms of formation of VO defect and its conversion to VO<sub>2</sub> defect.

Notably, the VO is optically active giving rise to an IR band at 830 cm<sup>-1</sup> in the neutral charge state and an IR band at 885 cm<sup>-1</sup> in the negative charge state [27,28]. Also, VO<sub>2</sub> defect gives an IR band at 888 cm<sup>-1</sup>, although it is electrically inactive [27,29]. Therefore IR spectroscopy is a suitable tool to make proper investigation on the properties and behavior of the above oxygen – vacancy defects.

## 2 Experimental methods

We used samples cut from prepolished Czochralski Si wafers. Their corresponding carbon and oxygen concentrations are given in Table I. Three groups of samples were used, irradiated with 2 MeV electrons at three fluences: group one (samples labeled S1i,  $i = 1-4$ ) with  $5 \times 10^{17} \text{ cm}^{-2}$ , group two (samples labeled S2i,  $i = 1-3$ ) with  $1 \times 10^{18} \text{ cm}^{-2}$  and group three (samples labeled S3i,  $i = 1-2$ ) with  $2 \times 10^{18} \text{ cm}^{-2}$ , using the Dynamitron accelerator at Takasaki-JAERI (Japan). In the labels of the samples index  $i$  increases with the increase of the carbon content of the sample. The samples S12 and S13 are doped with Ge concentrations of  $1 \times 10^{17}$  and  $1 \times 10^{18} \text{ cm}^{-3}$ , respectively. The sample S23 is doped with Sn concentration of  $3 \times 10^{17} \text{ cm}^{-3}$ . These Ge and Sn concentrations have practically a negligible effect [30-32] on the production and evolution of VO defect and for the purpose of this work are simply considered as Cz-Si samples containing carbon. After the irradiation, all the samples were subjected to isochronal anneals up to  $600 \text{ }^\circ\text{C}$ , in steps of  $\Delta T \approx 10 \text{ }^\circ\text{C}$  and  $\Delta t = 20 \text{ min}$ . After each annealing step, the IR spectra were recorded at room temperature by means of a FTIR spectrometer with a resolution of  $1 \text{ cm}^{-1}$ . The two phonon background absorption was subtracted from each spectrum by using a float-zone sample of equal thickness.

## 3 Results and discussion

Fig. 1 presents the IR spectrum of a representative sample S22 among the group of samples used in this experiment. It shows the spectrum after irradiation and at characteristic temperatures in the course of the isochronal anneal sequence. The above annealing temperatures were purposely chosen to focus on the production and annealing of the main irradiation-induced defects of VO ( $830 \text{ cm}^{-1}$ ) and  $\text{C}_i\text{O}_i$  ( $862 \text{ cm}^{-1}$ ) and their products upon annealing that is the  $\text{VO}_2$  ( $888 \text{ cm}^{-1}$ ),  $\text{VO}_3$  ( $904, 968, 1000 \text{ cm}^{-1}$ ) and the  $\text{C}_s\text{O}_{2i}$  ( $1048 \text{ cm}^{-1}$ ) [19] defects. In the spectra weak bands also appear and these are attributed [19] to the  $\text{C}_i\text{C}_s$

(546  $\text{cm}^{-1}$ ),  $\text{C}_s\text{O}_i$  (585, 637, 684  $\text{cm}^{-1}$ ) and the  $\text{C}_i\text{O}_i\text{Si}_I$  (936, 1020  $\text{cm}^{-1}$ ) complexes. Another band [19] of  $\text{C}_i\text{O}_i$  at 739  $\text{cm}^{-1}$  is also shown.

Fig. 2 shows the production of the VO defect as a function of the carbon concentration for the various irradiation fluences. Fig. 3 depicts how VO production varies with carbon content for each particular fluence used in the present study and in essence is an alternative representation of Fig. 2. As it is observed from Figs. 2 and 3 the concentration of VO defect increases with the increase of the carbon concentration. The relation is not linear, an indication that carbon affects indirectly the production of VO via the reactions taking place in the course of irradiation. Indeed, during irradiation V and  $\text{Si}_I$  are formed. Most of them are annihilated ( $V + \text{Si}_I \rightarrow 0$ ). The remaining vacancies are captured by oxygen atoms to form VO defects. The presence of carbon has the following effect: Due to its strong tendency to trap  $\text{Si}_I$ , carbon competes with vacancies in their capture. This leads in a decrease in the rate of the annihilation reaction ( $V + \text{Si}_I \rightarrow 0$ ), namely in an increase of the availability of vacancies to be captured by oxygen atoms. Thus the production of VO pairs increases, a fact verified experimentally (refer to Fig. 2). Additionally, it should be noted that in order to form a VO pair, a vacancy diffuses to an oxygen atom. It has been established experimentally that vacancies are mobile below room temperature. The onset of migration occurs [33] at  $\sim 70$  K in *n*-type Si, at  $\sim 150$  K in *p*-type Si and at  $\sim 200$  K in high resistivity material with corresponding migration activation energies 0.18 eV for *n*-type, 0.32 eV for *p*-type and 0.45 eV for the high resistivity Si. These energies are much lower than that of 2.53 eV for oxygen diffusion [13]. The disappearance of vacancies upon annealing is accompanied in the spectra by the emergence of other signals, identified as vacancies trapped by other defects [34]. This clearly confirms [34] that the annealing occurs as a long range migration of vacancies. In this process a vacancy performs a number of jumps in the lattice until it encounters an oxygen atom. This is considered as a long-distance migration process in the course of which the

vacancy travels through the lattice where it may encounter self-interstitials. In the presence of carbon this possibility is substantially decreased and the possibility of the traveling vacancy to encounter and pair with an oxygen atom is enhanced. In Fig. 3 we observe that with the increase of fluence in the range of  $5 \times 10^{17} - 2 \times 10^{18} \text{ cm}^{-2}$  the concentration of VO increases as expected [35,36]. Notably, there were reports in the past [37] that variations of carbon content have no impact on VO formation in contrast with these studies. These observations can be attributed to the quality of the Cz-Si samples which ~50 years ago was inferior.

Upon thermal anneal the  $830 \text{ cm}^{-1}$  VO band begins to disappear from the spectra and another band at  $888 \text{ cm}^{-1}$  attributed to the  $\text{VO}_2$  defect [27,29,38,39] begins to grow in the spectra. Fig. 4 shows the evolution of the VO and the  $\text{VO}_2$  bands for the S22 sample. It also shows the evolution of the  $\text{C}_i\text{O}_i$  and  $\text{C}_s\text{O}_{2i}$  bands to be discussed later. Fig. 5 shows the conversion ratio of the  $a_{\text{VO}_2}/a_{\text{VO}}$  versus carbon concentration for various fluencies. Fig. 6 is in essence an alternative representation of Fig. 5. It shows the conversion ratio of the  $a_{\text{VO}_2}/a_{\text{VO}}$  for various carbon concentrations, for each particular fluence used in this study. As it is immediately seen from the latter figures, the conversion ratio  $a_{\text{VO}_2}/a_{\text{VO}}$  decreases with the increase of the carbon concentration. To understand this behaviour one has to consider the main reactions that the VO participates upon annealing and how the formation of the  $\text{VO}_2$  defect is affected by the presence of carbon. Three main processes have been suggested in the literature [1,24,27,38,39] to contribute in the demise of the VO pair upon annealing. The first is a fast process related with the destruction of A-centers by self-interstitials:  $\text{VO} + \text{Si}_i \rightarrow \text{O}_i$  (1). The second is a slow process related with the migration of VO pair and its capture by immobile oxygen atoms to form the  $\text{VO}_2$  defect. Indeed the decrease of the ( $830 \text{ cm}^{-1}$ ) IR band of VO in the spectra upon anneal above  $300 \text{ }^\circ\text{C}$  is accompanied by the emergence of a band at ( $888 \text{ cm}^{-1}$ ) related [27] with the  $\text{VO}_2$  defect  $\text{VO} + \text{O}_i \rightarrow \text{VO}_2$  (2). However, annealing kinetics studies and comparisons of the concentration of the VO and the  $\text{VO}_2$  defects has led

to the suggestion [27,29] that a third process should take place involving the capture of VO **with** some unknown sinks, or/and the reaction of VO by other defects and impurities:  $VO + X \rightarrow VOX$  (3). It has been found [29] that X depends on the amount of carbon in Si material and it was suggested [29,39,40] that these unknown defects may be carbon related. Another relative suggested reaction is:  $VO + X \rightarrow X' + O_i$  (4). Also the following relative reactions have been considered [41] ( $VO + X \rightarrow XO + V$ ,  $VO + XO \rightarrow XVO_2$ ) (5). Furthermore, it deserves noting that reactions as VO dissociation ( $VO \rightarrow V + O_i$ ) (6), VO pairing with itself ( $VO + VO \rightarrow V_2O_2$ ,  $V_2O_2 \rightarrow VO_2 + V$ ) (7) and VO pairing with vacancies ( $VO + V \rightarrow V_2O$ ) (8) have also been considered and discussed in the literature [1,38,39]. It has been found [29] that more than half of the existed A-centers are not converted to  $VO_2$  defects. The question about the lack of the VO pairs **that are** not transformed to  $VO_2$  defects is still more or less unanswered and a complete picture about the fate of the loss VO pairs upon annealing has not been established so far. For instance, there are indications [40] **that** an electrical level  $E_c-0.20$  eV and a Photoluminescence signal at 950-meV may be related with a defect formed upon annealing of the VO pair but this needs further verification. Also the hypothesis that the VO pair is missing due to its capture by some infrared inactive sinks needs further information to **be verified**. In this work in order to understand the influence of carbon in the conversion of VO to  $VO_2$  we have firstly to take into account the reaction processes and especially those expressed by the relations (3), (4) and (5), indicating a direct or indirect role of carbon. For instance if X is carbon in reaction (3) then the formation of CVO defect is expected. Density functional theory (DFT) calculations indicate that the formation of the  $C_sVO$  defect is favourable (by -1.66 eV) [42]. This **illustrates** that the  $C_sVO$  defects are more bound as compared to the VO pairs (-1.32 eV). The DFT calculations also **show** that next nearest neighbor configurations of the  $C_sVO$  cluster are highly bound [42]. In essence when the VO pairs encounter a substitutional carbon atom at nearest neighbor or next nearest neighbor site,



they will be trapped and form the larger  $C_sVO$  cluster [42]. There are experimental reports [19] about the existence of a luminescence signal produced in irradiated Si, with a zero-phonon line at 488 meV ( $3942\text{ cm}^{-1}$ ) correlated with a  $(C_iO_i + V)$  structure. The suggested [19] reaction scheme for the production of the center is:  $C_s + Si_I \rightarrow C_i$ ,  $C_i + O_i \rightarrow C_iO_i$ ,  $(C_iO_i + V) \rightarrow 488\text{ meV band}$  and this defect anneals out at about  $200\text{ }^\circ\text{C}$ . These assignments are hypothetical and in view of the recent DFT evidence [42] it is possible to suggest an alternative reaction scheme ( $VO + C_s \rightarrow C_sVO$ ) leading to the formation of the  $C_sVO$  defects. In this scheme the  $200\text{ }^\circ\text{C}$  annealing temperature reflects the binding energy difference (between the  $C_sVO$  and  $VO$ , i.e.  $-0.34\text{ eV}$  when considering nearest neighbor  $C_sVO$  defects) required to dissociate the  $VO$  defect from the  $C_s$ .

If X is a carbon-related defect, a potential candidate is the  $C_iO_i$  complex. Indeed, the reaction  $VO + C_iO_i \rightarrow C_sO_{2i}$  has been reported previously [43]. The  $VO$  and  $C_iO_i$  bands anneal out together from the spectra at about  $300\text{ }^\circ\text{C}$  and a band at  $1048\text{ cm}^{-1}$  has been correlated with the  $C_sO_{2i}$  defect. The whole process is depicted in Fig. 4. However, the annealing of the  $C_iO_i$  defect [24] occurs mainly by dissociation so that a small percentage is responsible for the  $VO$  loss in the course of annealing, as a consequence of reactions with carbon-related complexes. Thus other complementary reactions are taking place. In this line of thought the role of carbon in the reduction of the ratio  $a_{VO_2}/a_{VO}$  could be mainly determined by the reaction  $C_s + Si_I \rightarrow C_i$  and in particular its effect on the reaction  $VO + O_i \rightarrow VO_2$ . The following approach is envisaged. On thermal anneals and for temperatures above  $300\text{ }^\circ\text{C}$ , carbon may capture only temporarily the self-interstitials. In essence, at these temperature carbon interstitials only temporarily form and then carbon atoms return back to substitutional sites releasing these self-interstitials ( $C_i \rightarrow C_s + Si_I$ ). A percentage of these self-interstitials then reacts with  $VO_2$  defects. In this sense, with the increase of carbon, the number of  $VO_2$  defects formed decreases. It is possible that some of the  $C_i$ 's could react with  $O_i$  atoms to

form  $C_iO_i$  pairs, something consistent with the slight increase of the latter defect prior to annealing that is in the temperature range 250-300 °C. Of course at these temperatures most of the  $C_i$  return at substitutional sites liberating  $Si_I$ 's, which react with  $VO_2$  defects as mentioned above. One may also consider that some of these self-interstitials react with  $VO$  defects ( $VO + Si_I \rightarrow O_i$ ), besides those reacting with  $VO_2$  defects. Thus, the number of  $VO$  pairs available to be converted to  $VO_2$  defects becomes also smaller. We argue that there is a balance between these two processes so that on the whole the presence of carbon leads to less  $VO_2$  defects in relation with the initial  $VO$  defects. In other words, less  $VO$  are converted to  $VO_2$  in the course of annealing, when carbon is present. This is equivalent with saying the ratio  $a_{VO_2}/a_{VO}$  decreases with the increase of the carbon concentration. In this framework it deserves mentioning that a defect reaction of  $VO_2$  defect with  $C_i$  has been suggested [44] to occur in high temperature (600-800 K) electron irradiations according to the reaction ( $VO_2 + C_i \rightarrow C_iVO_2$  or  $C_sO_{2i}$ ). Again in view of the recent DFT evidence [42] it is possible to suggest alternative reaction schemes ( $VO_2 + C_s \rightarrow C_sVO_2$  or  $O_i + C_sVO \rightarrow C_sVO_2$ ) leading to the formation of the  $C_sVO_2$  defects. The possibility of the existence of these extended clusters needs to be investigated further.

To explain the effect of carbon on the production of  $VO$  defect and its conversion to  $VO_2$  defect we have suggested a different role for carbon in the two cases. In the production process carbon traps self-interstitials thus suppressing the annihilation ratio with vacancies, which lead to an increase of the available vacancies leading to an enhancement of  $VO$  productivity. However, in the annealing process of  $VO$  and its conversion to the  $VO_2$  defect, carbon act as temporary trap for the self-interstitials which upon releasing they destroy some of the  $VO_2$  defects leading finally to a reduction of  $VO_2$  defects, alternatively to a diminishing of the  $a_{VO_2}/a_{VO}$  ratio. Thus, both processes, namely the production of  $VO$  defect and its

conversion to the  $VO_2$  defect are limited by reactions with self-interstitials the availability of which and their role is determined by the presence of carbon impurity.

Fig. 7 shows the conversion ratio of the  $a_{VO_3}/a_{VO_2}$  versus carbon concentration for various fluencies. Fig. 8 is in essence an alternative representation of Fig. 7. It shows the conversion ratio of the  $a_{VO_3} / a_{VO_2}$  for various carbon concentrations, for each particular fluence used in this study. It is observed that this conversion ratio between  $VO_2$  and  $VO_3$  defects decreases with the increase of the carbon concentration as opposed to the case between the  $VO_2$  and  $VO_3$  defects. Notably however, in the occasion of the  $VO_3$  defect its formation occurs at temperatures above  $450\text{ }^\circ\text{C}$  through the reaction  $VO_2 + O_i \rightarrow VO_3$ . At these temperatures oxygen impurity is already mobile [1], able to diffuse in the lattice and meet with the  $VO_2$  defect. In other words, in the case of the  $VO_3$  formation both partners that is  $VO_2$  defect [24] and oxygen atoms are moving in the lattice attaching each other in a  $VO_2 + O_i$  structure. Observations of an enhancement of oxygen diffusion in carbon containing Cz-Si, subjected to heat treatments at  $\sim 750\text{ }^\circ\text{C}$ , has been attributed [45] to the formation of fast diffusing oxygen –carbon molecules. We can extend this idea in the case of irradiated Si and for anneals above  $450\text{ }^\circ\text{C}$ . In particular we envisage the formation and rapid dissociation of oxygen –carbon complexes, a process which promotes in essence the diffusion of the oxygen impurity in the lattice. More specifically, it has been suggested [1] that large clusters formed as a result of the irradiation can release self-interstitials with the increase of temperature. In that case carbon interstitials form ( $C_s + Si_I \rightarrow C_i$ ), which are very mobile and they are captured by oxygen interstitials atoms to form  $C_iO_i$  defects. Above  $450\text{ }^\circ\text{C}$  these defects are unstable and they dissociate almost immediately only for the carbon atom to be trapped again by oxygen atom and then detrapping and so on. In other words, these oxygen-carbon complexes act as transient species and by their sequential formation and dissociation, in a self-sustaining loop at these temperatures, they could enhance the oxygen diffusion. Notably,

a similar behavior has been suggested [46] by the interaction of oxygen atoms with intrinsic defects. In this case the formation of oxygen-vacancy transient pairs could enhance oxygen diffusivity as well. Going one step forward in our picture, the enhanced oxygen diffusivity due to the presence of carbon could result in an increase of the probability of encountering  $VO_2$  defects with  $O_i$  atoms, as can be inferred by simple reaction kinetics arguments, leading finally to an increase of the total number of the formed  $VO_3$  defects, that is to an enhancement of the  $a_{VO_3} / a_{VO_2}$  ratio.

Interestingly, the role of carbon on the sequential formation of  $VO_n$  defects is different, depending on the temperature that determines the particular conversion. In the case of  $VO_2$  formation carbon affects reactions with self-interstitials although in the case of  $VO_3$  defects we have considered the likelihood of the formation of transient species involving carbon in their structures to understand the experimental observations. The present study is part of a concerted effort by the community to understand the impact of isovalent doping and/or localized strain on the defect processes of semiconductors [47-49].

#### **4 Conclusions**

We have investigated the effect of carbon on the properties of  $VO$  defect in electron-irradiated Cz-Si. In particular, we have focused our studies on the effect of carbon on the production  $VO$  defect and on its conversion to the  $VO_2$  defect, upon thermal annealing. It was found that when carbon concentration increases the production of  $VO$  defect is enhanced although its conversion to the  $VO_2$  defect is suppressed. The phenomena were attributed to the ability of carbon to trap self-interstitials thus affecting the balance among the reactions that govern the  $VO$  production and its conversion to the  $VO_2$  defect. Additionally, the  $a_{VO_3} / a_{VO_2}$  ratio increases with the increase of the carbon content as a result of the enhanced effect of carbon on the oxygen diffusivity of Si.

## References

1. R. C. Newman and R. Jones, in *Oxygen in Silicon*, edited by F. Shimura, Semiconductors and Semimetals Vol. 42 (Academic Press, Orlando, 1994), p. 289
2. C. Gao, X. Ma, J. Zhao, and D. Yang, *J. Appl. Phys.* **113**, 093511 (2013)
3. H. Wang, A. Chroneos, C. A. Londos, E. N. Sgourou, and U. Schwingenschlögl, *Appl. Phys. Lett.* **103**, 052101 (2013).
4. L. C. Kimerling, *Radiation Effects in Semiconductors*, Ed. N. B. Urli and J. W. Corbett, IOP Conf. Ser. No. 31, Bristol/ London 1977 (p.221).
5. C. A. Londos, *Phys. Stat. Solidi A* **113**, 503 (1989).
6. C. A. Londos, *Phys. Stat. Solidi A* **92**, 609 (1985).
7. S. D. Brotherton and P. Bradley, *J. Appl. Phys.* **53**, 5720 (1982).
8. A. Khan, M. Yamaguchi, Y. Ohshita, N. Dharmarasu, K. Araki, T. Abe, H. Itoh, T. Ohshima, M. Imaizumi, and S. Matsuda, *J. Appl. Phys.* **90**, 1170 (2001).
9. C. A. Londos, L. G. Fytros, and G. J. Georgiou, *Defect and Diffusion Forum*, **171-172**, 1 (1999).
10. K. Gill, G. Hall, and B. MacEvoy, *J. Appl. Phys.* **82**, 126 (1997).
11. B. G. Svensson and J. L. Lindstrom and J. W. Corbett, *Appl. Phys. Lett.* **47**, 841 (1985).
12. R. C. Newman, *J. Phys.: Condens. Matter* **12**, R335 (2000).
13. J. C. Mikkelsen, in *Oxygen, Carbon, Hydrogen, and Nitrogen in Crystalline Silicon*, MRS Symposia Proceedings No 59, edited by J. C. Mikkelsen, Jr., S. J. Pearton, J. W. Corbett, and P. W. Pennycook (Materials Research Society, Pittsburgh, 1986), p. 19
14. P. Wagner, J. Hage, J. M. Trombetta, and G. D. Watkins, *Mater. Sci. Forum* **83-87**, 401 (1992).
15. L. I. Murin, V.P. Markevich, I. F. Medvedeva, and L. Dobaczewski, *Semiconductors* **40**, 1282 (2006).
16. C. A. Londos, N. Sarlis, L. G. Fytros and K. Papastergiou, *Phys. Rev. B* **53**, 6900 (1996)
17. V. V. Voronkov and R. Falster, *J. Electrochem. Soc.* **149**, G167 (2002).
18. G. Kissinger, J. Dadrowski, A. Sattler, C. Seuring, T. Muller, H. Richter, W. Von Ammon, *J. Electrochem. Soc.* **154**, H454 (2007).

19. G. Davies and R. C. Newman, in *Handbook on Semiconductors, Materials Properties and Preparations*, edited by T.S. Moss and S. Mahajan (North Holland, Amsterdam 1994), p. 1557
20. S. G. Cloutier, P. A. Kossyrev, and J. Xu, *Nat. Mater.* **4**, 877 (2005)
21. E. Rotem, J. M. Shainline, and J. M. Xu, *Appl. Phys. Lett.* **91**, 051127 (2007)
22. D. D. Berhanuddin, M. A. Lourenço, R. M. Gwilliam, and K. P. Homewood, *Adv. Funct. Mater.* **22**, 2709 (2012)
23. C. A. Londos, A. Andrianakis, V. Emtsev, and H. Ohyama, *J. Appl. Phys.* **105**, 123508 (2009).
24. B. G. Svensson and J. L. Lindstrom, *phys. stat. sol. (a)* **95**, 537 (1986).
25. C. A. Londos, *Jpn J. Appl. Phys., part I* **27**, 2089 (1988).
26. C. A. Londos, *Phys. Stat. Solidi A* **102**, 639 (1987)
27. J. W. Corbett., G. D. Watkins, and R. S. McDonald *Phys. Rev.* **135**, A1381 (1964).
28. A. R. Bean and R. C. Newman *Solid State Commun.* **9**, 271 (1971).
29. J. L. Lindstrom and B. G. Svensson, in *Oxygen, Carbon, Hydrogen, and Nitrogen in Crystalline Silicon*, MRS Symposia Proceedings No 59, edited by J. C. Mikkelsen, Jr., S. J. Pearton, J. W. Corbett, and P. W. Pennycook (Materials Research Society, Pittsburgh, 1986), p. 45
30. C. A. Londos, A. Andrianakis, V. V. Emtsev and H. Ohyama, *Semicond. Sci. Technol.* **24**, 075002 (2009).
31. A. Chroneos, C. A. Londos, E. N. Sgourou, and P. Pochet, *Appl. Phys. Lett.* **99**, 241901 (2011).
32. A. Chroneos, C. A. Londos, E. N. Sgourou, *J. Appl. Phys.* **110**, 093507 (2011).
33. G. D. Watkins, *Mater. Sci. Semicond. Proc.* **3**, 227 (2000).
34. G. D. Watkins, *Symposium of Radiation effects on Semiconductors, 1967*, (Toulouse, Journées d'Electronique), pA1-1.
35. J. W. Corbett and G. D. Watkins, *Phys. Rev.* **138**, A555 (1965).
36. J. W. Corbett, G. D. Watkins, R. M. Chrenko and R. S. McDoland, *Phys. Rev.* **121**, 1015 (1961).
37. A. R. Bean, R. C. Newman and R. S. Smith *J. Phys. Chem. Solids* **31**, 739 (1970).
38. C. A. Londos, N. V. Sarlis, and L. G. Fytros, *Phys. Stat. Solidi A* **163**, 325 (1997).

39. B. G. Svensson and J. L. Lindstrom, Phys. Rev. B **34**, 8709 (1986).
40. O. O. Awadelkarim, H. Weman, B. G. Svensson and J. L. Lindstrom, J. Appl. Phys. **60**, 1974 (1986).
41. L. F. Makarenko, Semicond. Sci. Technol. **8**, 1692 (1993).
42. A. Chroneos and C. A. Londos, J. Appl. Phys. **107**, 093518 (2010).
43. N. Inoue, H. Ohyama, Y. Goto, and T. Suriyama, Physica B **401-402**, 477 (2007).
44. J. L. Lindstrom, L. I. Murin, T. Hallberg, V. P. Markevich, B. G. Svensson, M. Kleverman, and J. Hermansson, Nucl. Instr. Meth. Phys. Res. B **186**, 121 (2002)
45. F. Shimura and T. Higuchi, R. S. Hockett, Appl. Phys. Lett. **53**, 69 (1988)
46. A. S. Oates and R. C. Newman, Appl. Phys. Lett. **49**, 262 (1986)
47. C. N. Koumelis, G. E. Zardas, C. A. Londos, and D. K. Leventouri, Acta Crystallogr. A **32**, 306 (1976).
48. A. Chroneos, H. Bracht, R. W. Grimes, and B. P. Uberuaga, Mater. Sci. Eng. B **154**, 72 (2008).
49. A. Chroneos, E. N. Sgourou, C. A. Londos, J. Mater. Sci.: Mater. Electron. **24**,2772 (2013).

**Table I** The concentration of VO and the ratio  $a_{VO_2}/a_{VO}$  and  $a_{VO_3}/a_{VO_2}$  for the used samples irradiated with 2MeV electrons at various fluences.

Sample New name	Sample Old name	$[C_s]_o$ $10^{16}$ ( $cm^{-3}$ )	$[C_s]_{a.i.}$ $10^{16}$ ( $cm^{-3}$ )	$[O_i]_o$ $10^{17}$ ( $cm^{-3}$ )	$[O_i]_{a.i.}$ $10^{17}$ ( $cm^{-3}$ )	Fluence ( $e\ cm^2$ )	$[VO]$ $10^{16}$	$a_{VO_2}/a_{VO}$	$a_{VO_3}/a_{VO_2}$
S11	E4	<1	-	9.6	9.2	$5 \times 10^{17}$	2.3	0.83	0.4
S12	Ge1	2	<2	9.6	9.05	$5 \times 10^{17}$	3.4	0.63	0.44
S13	Ge3	3	2	10	9.42	$5 \times 10^{17}$	3.5	0.62	0.48
S14	E3	16	9.6	10.2	9.8	$5 \times 10^{17}$	3.9	0.56	0.54
S21	M4	5	<2	9.5	9.04	$1 \times 10^{18}$	5.87	0.48	0.43
S22	H3	22	13.2	9.3	9	$1 \times 10^{18}$	6.2	0.41	0.46
S23	Sn3.1	47	18	9.6	8.5	$1 \times 10^{18}$	7.3	0.38	0.58
S31	M7	4	2.3	10.5	9.72	$2 \times 10^{18}$	7.2	0.55	0.33
S32	M5	8.8	3.85	9.45	8.69	$2 \times 10^{18}$	9.8	0.33	0.44



## Figure captions

Fig. 1. IR spectra of the S22 sample after irradiation and at 250, 400 and 500 °C in the course of anneals.

Fig. 2 The production of VO defect as a function of carbon concentration for the various fluences.

Fig. 3 The production of VO defect for the various carbon concentrations of the samples, for each particular fluence used in this work.

Fig. 4 Evolution of the IR bands of VO ( $830\text{ cm}^{-1}$ ),  $\text{C}_i\text{O}_i$  ( $862\text{ cm}^{-1}$ ),  $\text{VO}_2$  ( $888\text{ cm}^{-1}$ ) and  $\text{C}_s\text{O}_{2i}$  ( $1048\text{ cm}^{-3}$ ) defects in the ( $2\text{ MeV} / 1 \times 10^{18}\text{ cm}^{-2}$ ) electron-irradiated S22 sample upon isochronal annealing.

Fig. 5 The conversion ratio  $a_{\text{VO}_2}/a_{\text{VO}}$  as a function carbon concentration for the various fluences.

Fig. 6 The conversion ratio  $a_{\text{VO}_2}/a_{\text{VO}}$  for the various carbon concentrations, for each particular fluence used in this work.

Fig. 7 The conversion ratio  $a_{\text{VO}_3}/a_{\text{VO}_2}$  as a function carbon concentration for the various fluences.

Fig. 8 The conversion ratio  $a_{\text{VO}_3}/a_{\text{VO}_2}$  for the various carbon concentrations, for each particular fluence used in this work.

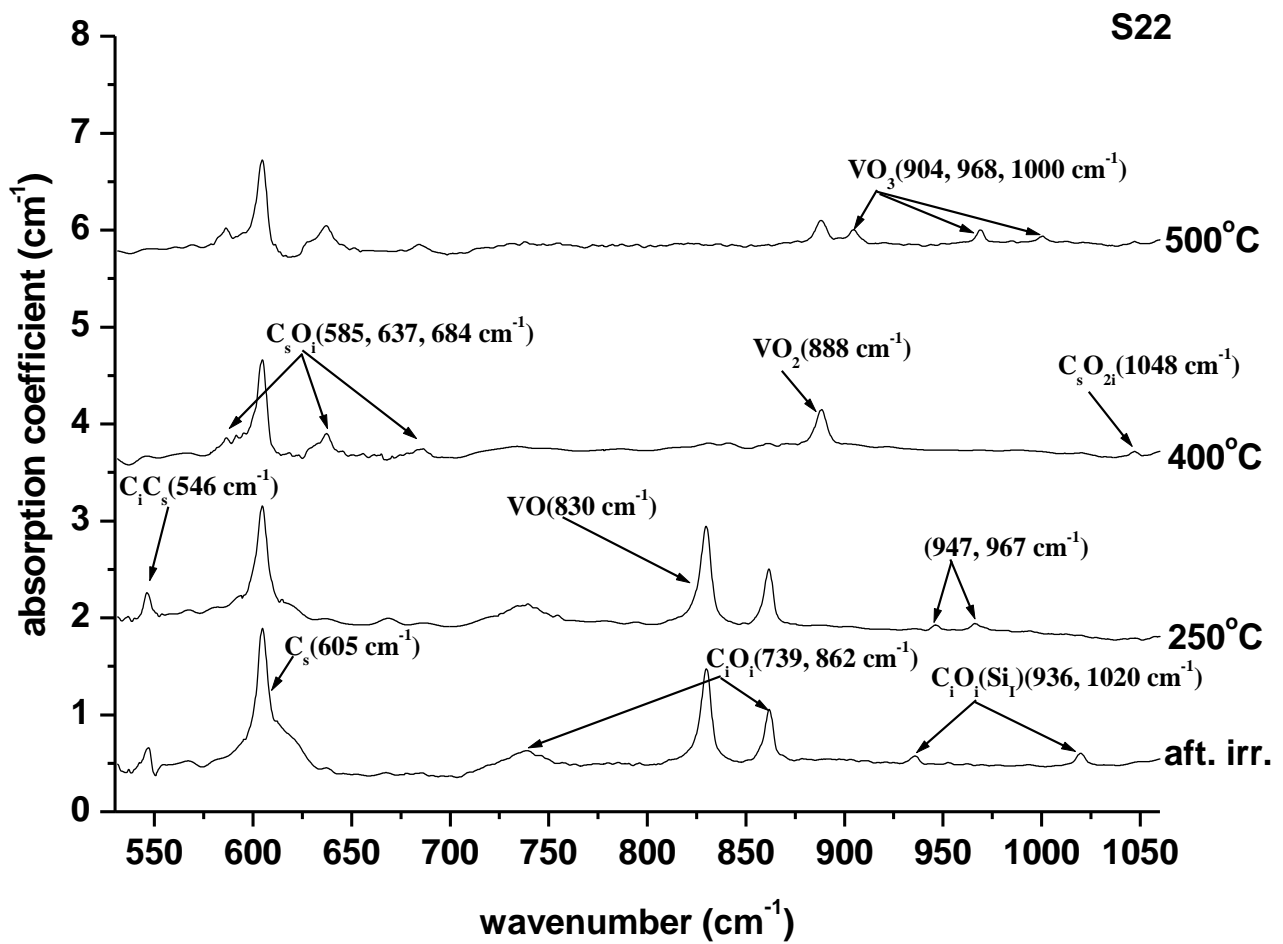


Fig.1

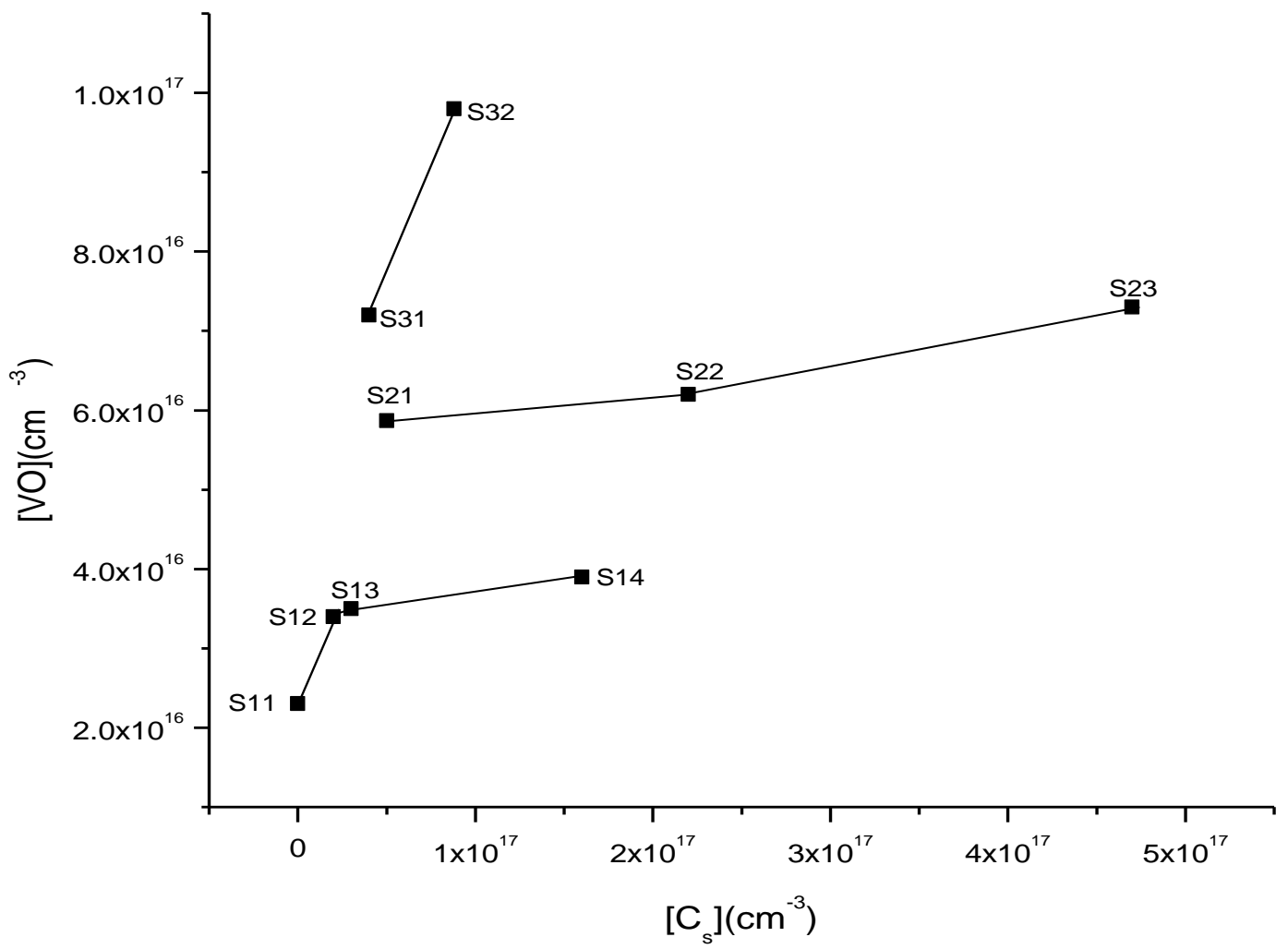


Fig.2

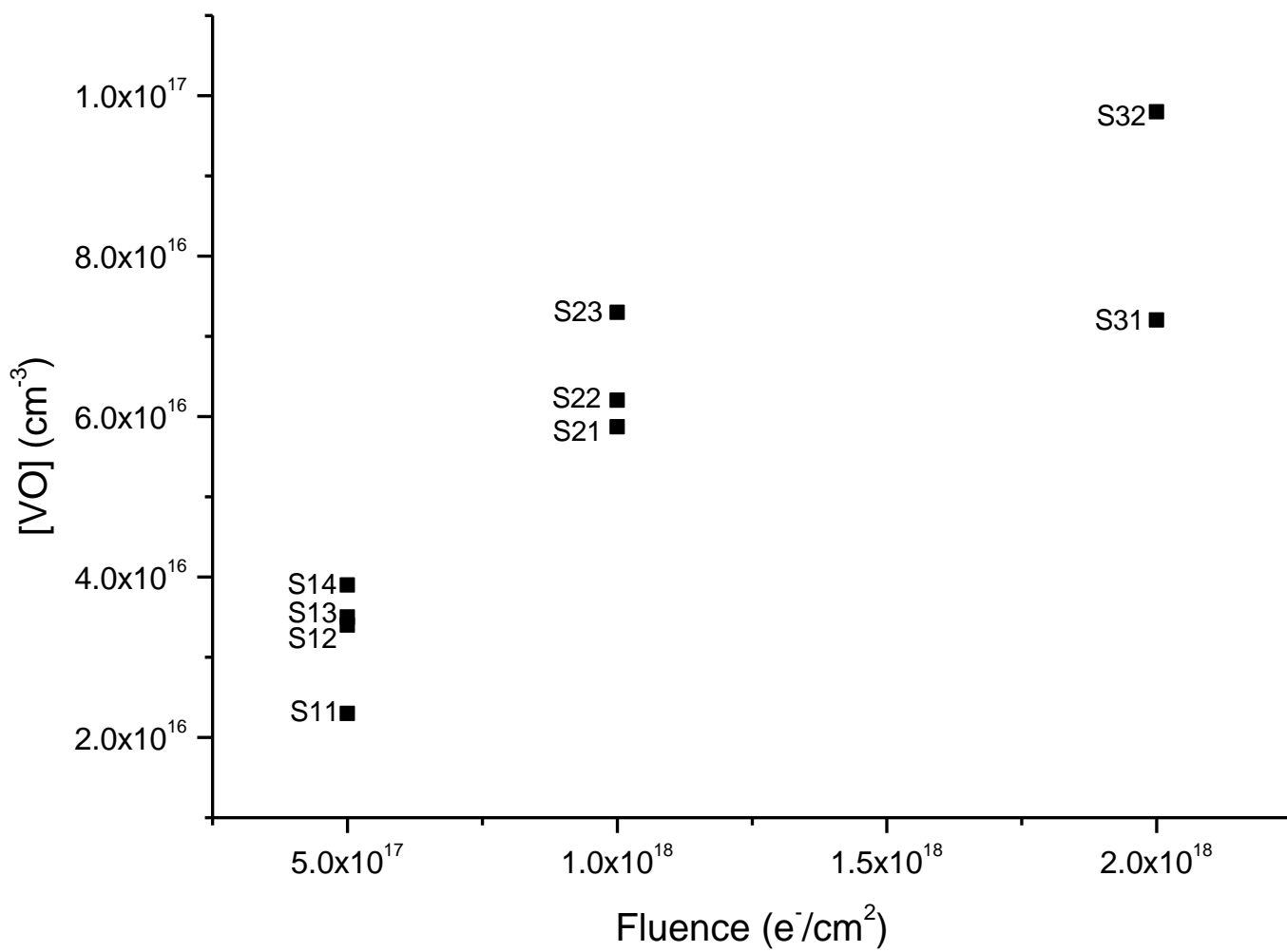


Fig. 3

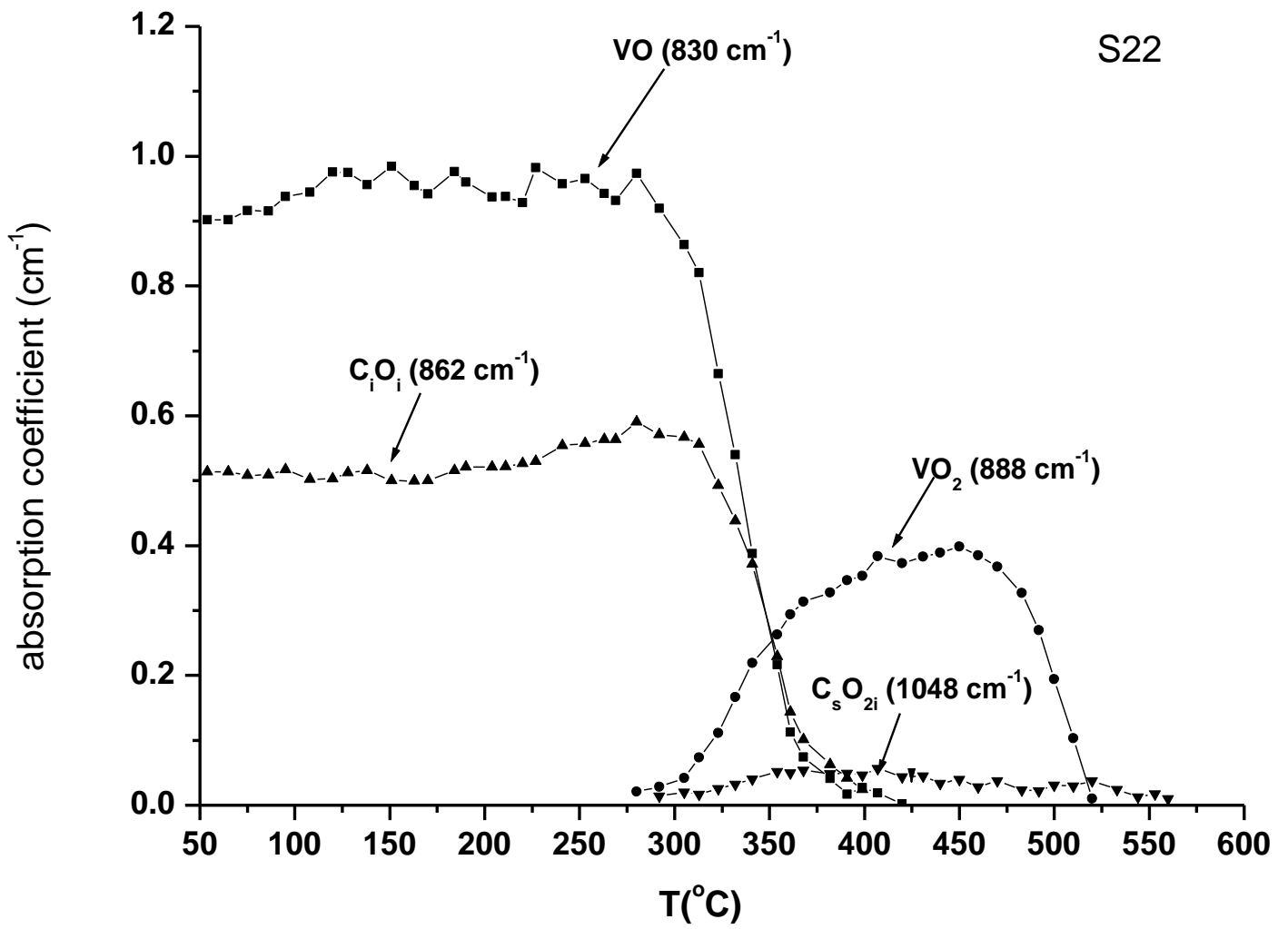


Fig.4

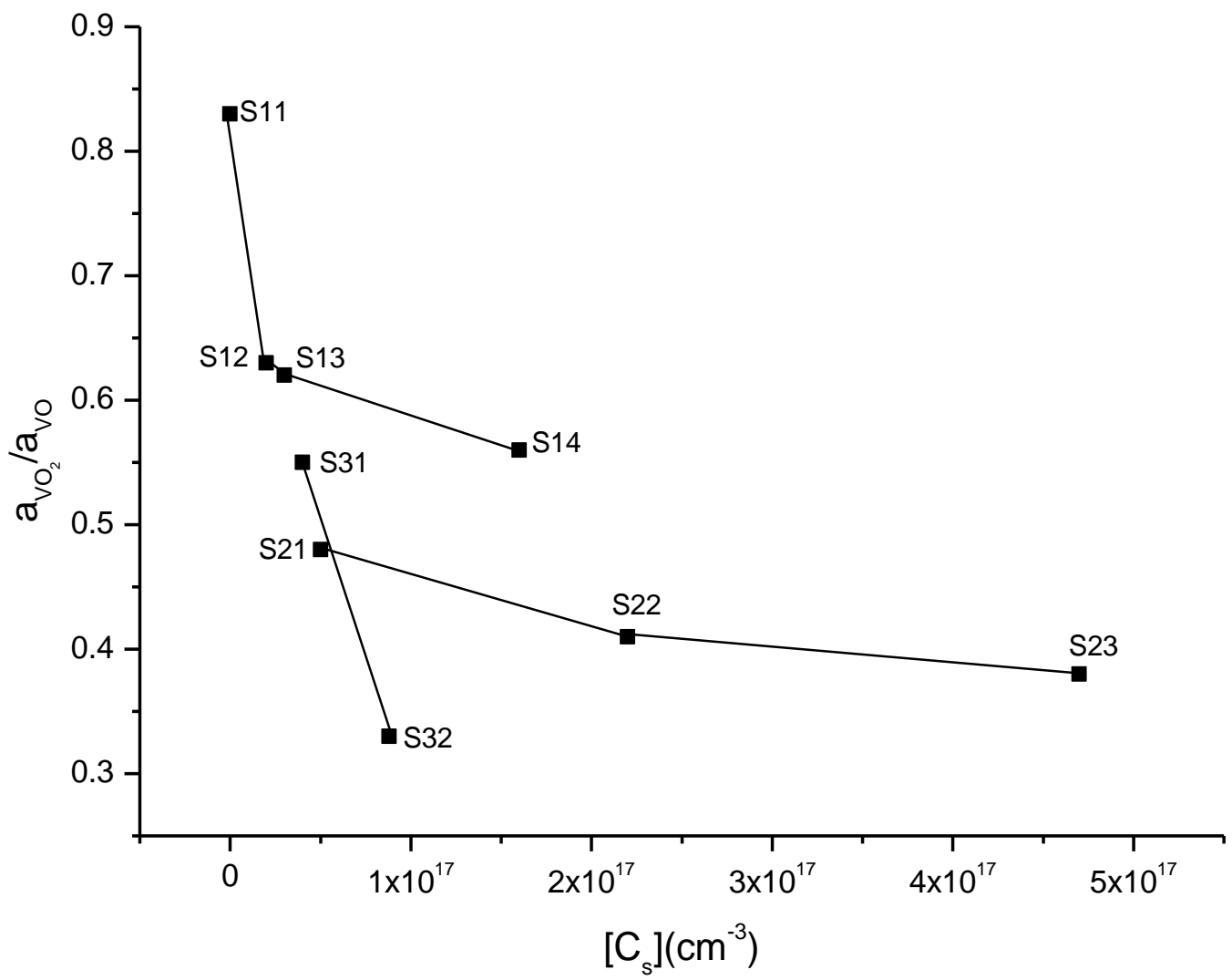


Fig.5

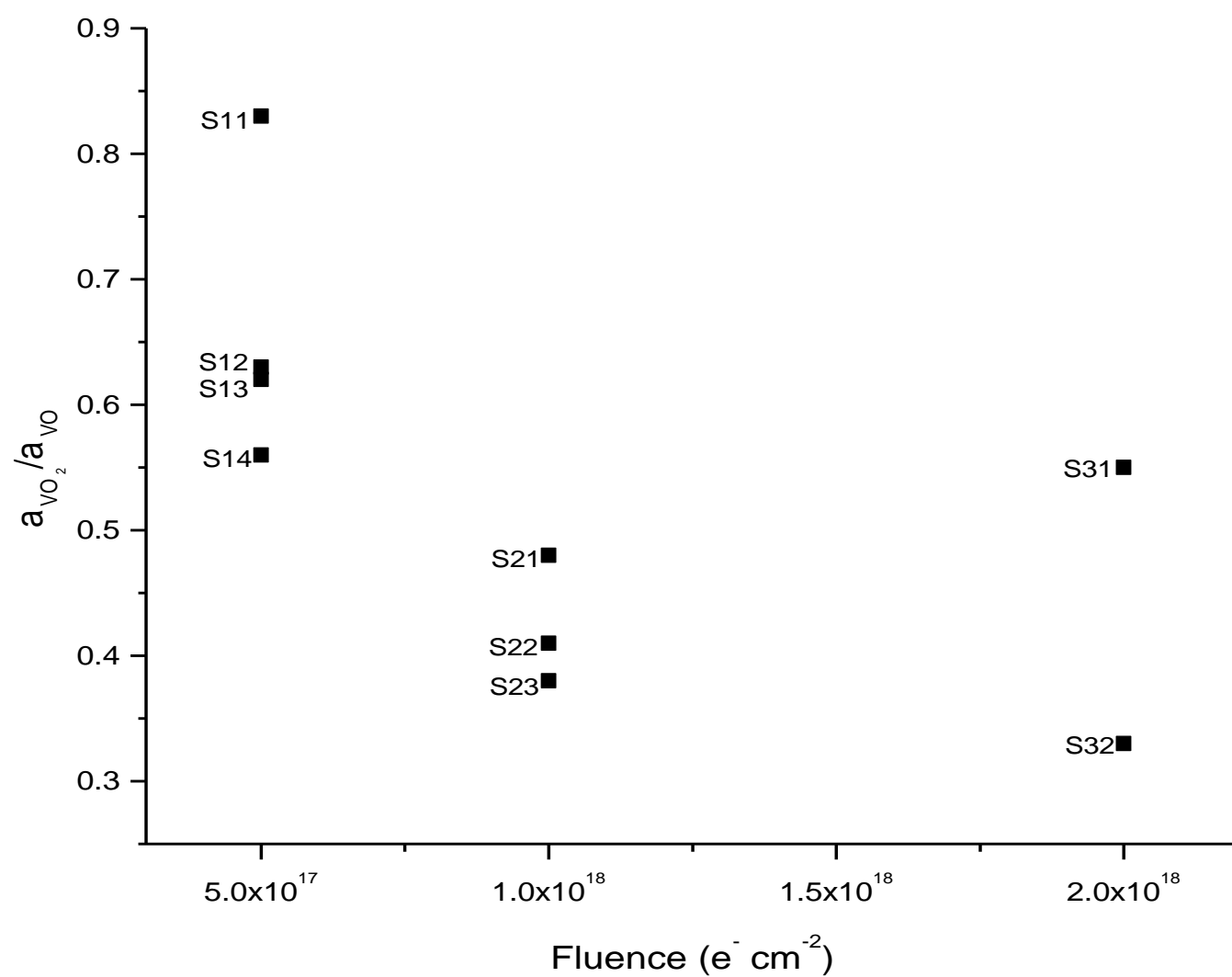


Fig. 6

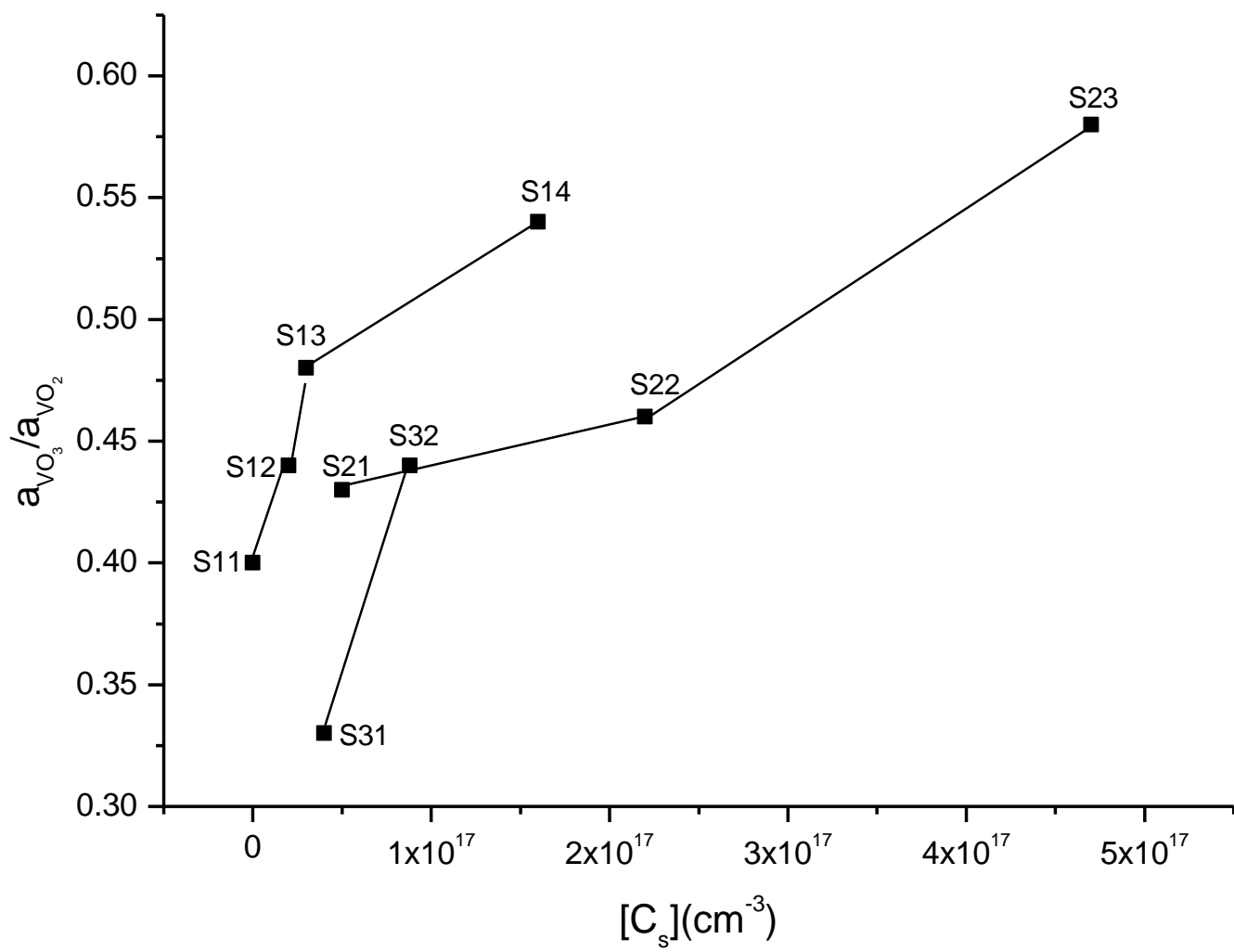
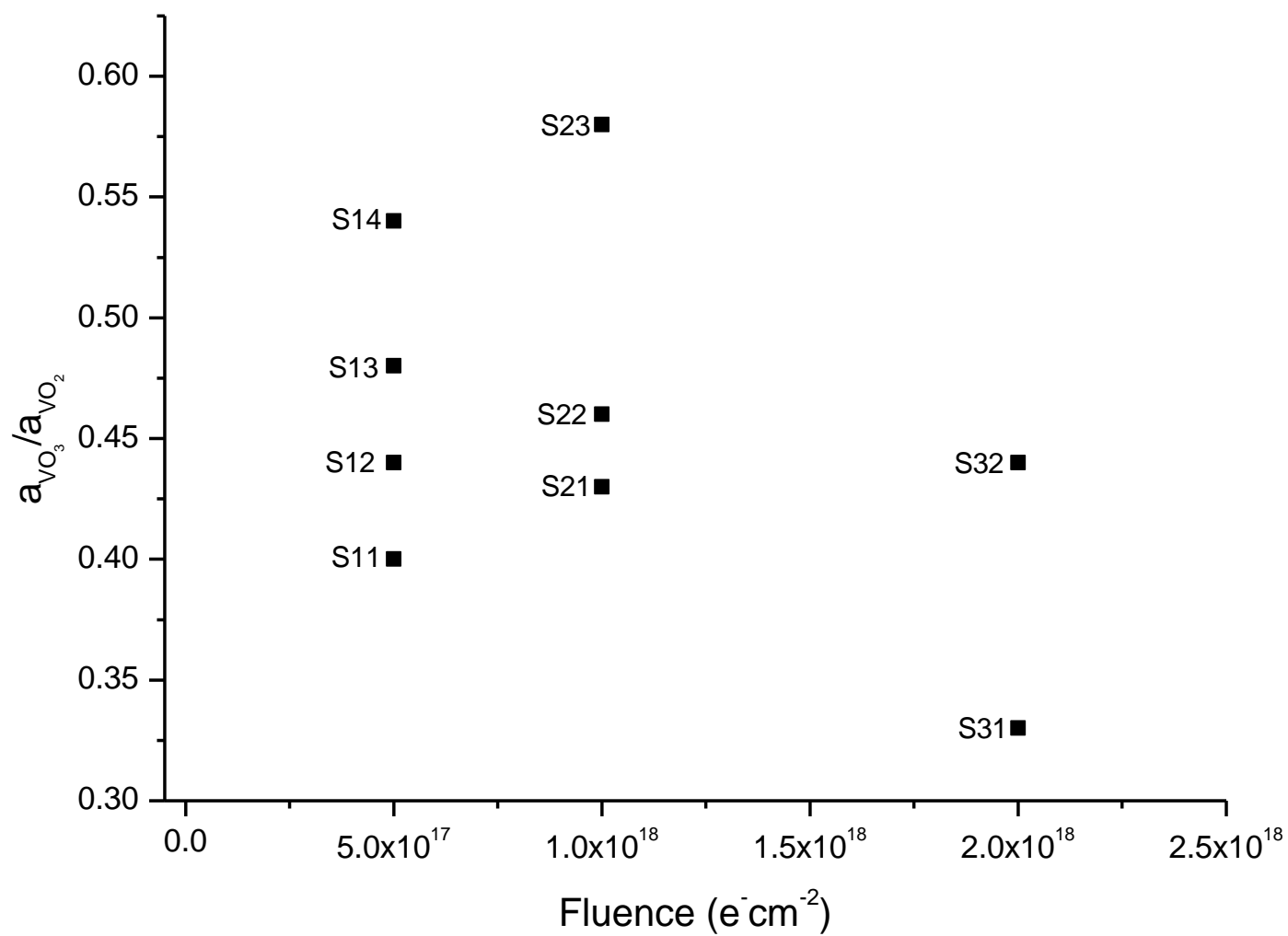


Fig.7





**Fig. 8**